

Remarks on “ COLORING RANDOM TRIANGULATION”

S.BALASKA¹ and W. RÜHL²

*Department of Physics, University of Kaiserslautern, P.O.Box 3049
67653 Kaiserslautern, Germany*

Abstract

We transform the two-matrix model, studied by P.Di Francesco and al. in [1], into a normal one-matrix model by identifying a “formal” integral used by these authors as a proper integral. We show also, using their method, that the results obtained for the resolvent and the density are not reliable.

1 THE MODEL

In a recent paper P. Di Francesco, B. Eynard and E. Gittert [1] discuss a model of two $n \times n$ hermitean matrices M and R with a partition function

$$Z(p, q, g; N) = \int dM dR \exp\{-N\text{Tr}[p \log(1 - M) + q \log(1 - R) + gMR]\} \quad (1)$$

where g is later replaced by

$$g = \frac{1}{t} \quad (2)$$

and the strong coupling limit is of primary interest: $t \rightarrow 0$. Applying the Itzykson-Zuber integral identity, the partition function is transformed into an integral over the

¹On leave of absence from Laboratoire de Physique Théorique, Université d’Oran, 31100 Es-Sénia, Algeria, E-mail: balaska@physik.uni-kl.de

²E-mail:ruehl@physik.uni-kl.de

eigenvalues $\{m_i\}_{i=1,\dots,n}$ and $\{r_i\}_{i=1,\dots,n}$ of M and R respectively.

$$Z(p, q, g; N) \sim \int \Delta(m) \Delta(r) \prod_{i=1}^n [\exp\{-N[p \log(1 - m_i) + q \log(1 - r_i) + gm_i r_i]\} dr_i dm_i] \quad (3)$$

where Δ is the Vandermonde determinant.

However, an integral such as

$$\int_{\mathbb{R}^2} dx dy \exp\{-N[p \log(1 - x) + q \log(1 - y) + gxy]\} p_n(x) \tilde{p}_m(y) \quad (4)$$

where $p_n(\tilde{p}_m)$ are polynomials of degree $n(m)$ does not exist.

In later parts of the paper the authors want to ascribe a meaning to such integrals by the “formal integral” (see eq. (4.2))

$$\int_{\mathbb{R}^2} dx dy x^\alpha y^\beta \exp\left[-\frac{N}{t} xy\right] = \alpha! \delta_{\alpha\beta} \left(\frac{t}{N}\right)^{\alpha+1} \quad (5)$$

$(\alpha, \beta \in \mathbb{Z}_+)$

Special cases of such “formal integral” are already used earlier in the text. It is, however, easy to see that the “formal integral” is a proper integral

$$2 \int_{\mathbb{R}^2} dx dy w^\alpha \bar{w}^\beta \exp\left[-\frac{N}{t} w \bar{w}\right] = \alpha! \delta_{\alpha\beta} \left(\frac{t}{N}\right)^{\alpha+1} \quad (6)$$

$(w = x + iy, \bar{w} = x - iy)$

It is also easy to go back and replace all “formal integrals” by proper integrals, then for (3) we obtain

$$Z(p, q, g; N) \sim \int |\Delta(w)|^2 \prod_{i=1}^n [\exp\{-N[p \log(1 - w_i) + q \log(1 - \bar{w}_i) + gw_i \bar{w}_i]\}] dx_i dy_i \quad (7)$$

where we have to restrict the w_i to, say

$$0 \leq |w_i| \leq 1 \quad (8)$$

and

$$\operatorname{Re}(Np) < 1, \operatorname{Re}(Nq) < 1 \quad (9)$$

if we want to give an analytic meaning to Z in the variable $t = g^{-1}$, but we can integrate over the whole complex plane if we are interested only in the formal power expansion in t for $t \rightarrow 0$. Finally, without referring to the Itzykson-Zuber formula we can integrate over $U(n)$ to get

$$Z(p, q, g; N) = 2^n \int dM dM^\dagger \exp\{-N \operatorname{Tr}[p \log(1 - M) + q \log(1 - M^\dagger) + g M M^\dagger]\} \quad (10)$$

where M is a normal matrix

$$[M, M^\dagger] = 0 \quad (11)$$

Thus the model actually evaluated is not a hermitean two-matrix model but a normal one-matrix model.

2 THE SADDLE POINT EQUATION

The partition function is calculated in the limit $N \rightarrow \infty$ so that $n = zN$ and z is kept fixed. In this limit the partition function is approximated by (see eq.(4.20) of [1])

$$Z = N^n \int_0^\infty d\alpha_i \dots d\alpha_n \exp\{N^2 S(\alpha_1 \dots \alpha_n, p, q, z)\} \quad (12)$$

with

$$\begin{aligned} S = \frac{1}{N} \sum_{i=1}^n & \left[(\alpha_i + p - z)(\log(\alpha_i + p - z) - 1) \right. \\ & + (\alpha_i + q - z)(\log(\alpha_i + q - z) - 1) - 2\alpha_i(\log(\alpha_i) - 1) \\ & + \frac{1}{N} \log(\gamma(\alpha_i N + 1, \frac{N}{t})) + \alpha_i \log(\frac{t}{N}) \left. \right] \\ & + \frac{1}{N^2} \sum_{i \neq j} \log(|\alpha_i - \alpha_j|) \\ & - \int_0^z ds [(p - s)(\log(p - s) - 1) + (q - s)(\log(q - s) - 1) \\ & \quad + s(\log(s) - 1) + s \log(t)] \end{aligned} \quad (13)$$

where $\gamma(\alpha, \xi)$ is the incomplete γ -function [2]. In the limit $N \rightarrow \infty$ the term containing the incomplete γ -function is calculated as

$$\lim_{N \rightarrow \infty} \left\{ -\frac{t}{N} \log(\gamma(\xi \frac{N}{t} + 1, \frac{N}{t})) - \xi \log(\frac{t}{N}) \right\} = \Theta(\xi - 1) + \xi(1 - \log(\xi))\Theta(1 - \xi) \quad (14)$$

where $\xi = \alpha t$ and so that the derivative of the r.h.s is continuous at $\xi = 1$. The saddle point equation is then

$$\begin{aligned} 2P \int d\beta \frac{\rho(\beta)}{\beta - \alpha} &= \Theta(1 - \alpha t) \log\left(t \frac{(\alpha + p - z)(\alpha + q - z)}{\alpha}\right) \\ &+ \Theta(\alpha t - 1) \log\left(\frac{(\alpha + p - z)(\alpha + q - z)}{\alpha^2}\right) \end{aligned} \quad (15)$$

valid for α in the support of the density ρ and with the normalisation

$$\int d\beta \rho(\beta) = z \quad (16)$$

We mention that the free energy also has an additional term compared with the expression obtained in the equation (4.29) of [1]

$$t \partial_t f = \int_0^\infty d\beta \beta \rho(\beta) - \frac{1}{2} z^2 + \frac{1}{t} \int_{t^{-1}}^\infty d\beta \rho(\beta) \quad (17)$$

We solve the saddle point equation by making the ansatz

$$\rho(\alpha) = \rho_1(\alpha) + \rho_2(\alpha) \quad (18)$$

$$\rho_i(\alpha) \geq 0 \quad , \quad \int d\beta \rho_i(\beta) = z_i \quad (19)$$

$$z_1 + z_2 = z \quad (20)$$

and

$$supp(\rho_1) = \langle \gamma_1, \gamma_2 \rangle \quad (21)$$

$$supp(\rho_2) = \langle \gamma_3, \gamma_4 \rangle \quad (22)$$

In order to solve (15) we have to make sure that

$$\gamma_2 \leq \frac{1}{t} \leq \gamma_3 \quad (23)$$

Next we introduce the resolvent

$$\omega_i(\alpha) = \int d\beta \frac{\rho_i(\beta)}{\alpha - \beta} \quad (24)$$

so that (15) induces

$$\omega_1(\alpha + i0) + \omega_1(\alpha - i0) = -\log\left(t \frac{(\alpha + p - z)(\alpha + q - z)}{\alpha}\right) \quad (25)$$

$$\omega_2(\alpha + i0) + \omega_2(\alpha - i0) = -\log\left(\frac{(\alpha + p - z)(\alpha + q - z)}{\alpha^2}\right) \quad (26)$$

and

$$\lim_{\alpha \rightarrow \infty} \alpha \omega_i(\alpha) = \int d\beta \rho_i(\beta) = z_i \quad (27)$$

We shall use the method of [1], and show that their result for ω_1, ρ_1 is not reliable, in spite of the fact that the saddle point equation (25) is the same as the one obtained in [1]. For simplicity we set $z_1 = z, z_2 = 0$ from now on.

We introduce two parameters r, δ and three hyperbolic angles (assumed to be all ≥ 0) Φ_1, Φ_2, Φ_3 by

$$\alpha = z - r - 2\delta \cosh(\Phi) \quad (28)$$

$$p = r + 2\delta \cosh(\Phi_1) \quad (29)$$

$$q = r + 2\delta \cosh(\Phi_2) \quad (30)$$

$$z = r + 2\delta \cosh(\Phi_3) \quad (31)$$

Then

$$\frac{t}{\alpha}(\alpha + p - z)(\alpha + q - z) = \delta t \frac{4(T_1^2 - T^2)(T_2^2 - T^2)(1 - T_3^2)}{(1 - T_1^2)(1 - T_2^2)(1 - T^2)(T_3^2 - T^2)} \quad (32)$$

where

$$T = \tanh\left(\frac{\Phi}{2}\right) \quad (33)$$

$$T_i = \tanh\left(\frac{\Phi_i}{2}\right) \geq 0 \quad (34)$$

with

$$\gamma_1 = z - r - 2\delta, \quad \gamma_2 = z - r + 2\delta$$

Along the cut $\langle \gamma_1, \gamma_2 \rangle$ the parameter Φ is assumed to vary as $+i\varphi$ ($0 \leq \varphi \leq \pi$) for $Im(\alpha) \searrow 0$ and as $-i\varphi$ for $Im(\alpha) \nearrow 0$. On the real axis outside the cut, Φ is necessarily so that

$$T = \tanh\left(\frac{\Phi}{2}\right) \rightarrow -1 \quad \text{for } \alpha \rightarrow \pm\infty \quad (35)$$

Taking into account (34) and (35) only one ansatz for $\omega_1(\alpha)$ is possible which has only the cut at $\langle \gamma_1, \gamma_2 \rangle$, namely

$$\omega_1(\alpha) = -\log\left(\frac{2(T_1 - T)(T_2 - T)(1 + T_3)}{(1 - T)(1 + T_1)(1 + T_2)(T_3 - T)}\right) \quad (36)$$

provided we take

$$\begin{aligned} \delta t &= \frac{(1 - T_1)(1 - T_2)(1 + T_3)}{(1 + T_1)(1 + T_2)(1 - T_3)} \\ &= u_1 u_2 u_3^{-1} \end{aligned} \quad (37)$$

where

$$\begin{aligned} u_i &= \frac{1 - T_i}{1 + T_i} = e^{-\Phi_i} \\ u &= \frac{1 - T}{1 + T} = e^{-\Phi} \end{aligned}$$

Finally we obtain for the density

$$\rho_1(\alpha) = \frac{1}{2\pi} \left(\varphi - 2(\Psi_1(\varphi) + \Psi_2(\varphi) - \Psi_3(\varphi)) \right) \quad (38)$$

where

$$\Psi(\alpha) = \arctan\left(\frac{\tan(\frac{\varphi}{2})}{T_i}\right)$$

We can show that

$$\begin{aligned} z &= \int \rho_1(\beta) d\beta \\ &= \int_0^\pi d\varphi 2\delta \sin(\varphi) \rho_1(\beta(\varphi)) \\ &= \delta[u_3 - u_1 - u_2] \\ &= \lim_{\alpha \rightarrow +\infty} \alpha \omega_1(\alpha) \end{aligned} \quad (39)$$

as desired.

Positivity of the density function (38) is achieved if and only if

$$T_1 + T_2 - T_3 - 1 \geq 0 \quad (40)$$

This constraint follows from the observation that the neighborhood of π in the variable φ is critical, i.e. negative values of ρ_1 appear here first. Namely, if we write

$$\frac{\varphi}{2} = \frac{\pi}{2} - \chi \quad (41)$$

we have then

$$\Psi_i(\varphi) = \frac{\pi}{2} - T_i \chi + O(\chi^3) \quad (42)$$

and consequently

$$\rho_1(\alpha) = \frac{1}{\pi} (T_1 + T_2 - T_3 - 1) \chi + O(\chi^3) \quad (43)$$

References

- [1] P. Di Francesco, B. Eynard and E. Guitter, *Coloring Random Triangulation* , cond-mat/9711050
- [2] M.Abramowitz and I.A.Stegun, *Handbook of Mathematical Functions*, Dover publications, INC., New York.